

PHYS 1211 F2020

Lecture Wednesday November  
4, 2020

Chapter 7

Introduction to Work by spring and variable force

# Calculus for 1<sup>st</sup> year physics: Derivatives and Indefinite Integrals

- Consider a function  $y$  of a variable  $x$ ,  $y(x)$
- **Derivative** of  $y$  wrt  $x$ ,  $\frac{dy}{dx} = y'$
- Example,  $y = x^2 \rightarrow \frac{dy}{dx} = \frac{d(x^2)}{dx} = 2 \times x^{2-1} = 2x$
- **Indefinite integral**,  $\int y \, dx = z(x)$ , where  $y(x)$  and  $z(x)$  are functions of  $x$ .
- The function  $z(x)$  is the **anti-derivative** of the function  $y(x)$ :  $\frac{dz}{dx} = y(x)$
- Example  $y = x \rightarrow \int y \, dx = \int x \, dx = \frac{1}{2}x^2 = z(x)$
- Since if  $z(x) = \frac{1}{2}x^2 \rightarrow \frac{dz}{dx} = \frac{d(\frac{1}{2}x^2)}{dx} = \frac{1}{2}2 \times x^{2-1} = x$

# Calculus for 1<sup>st</sup> year physics: Definite Integrals

- **Definite integral**  $\int_{x_i}^{x_f} y(x) dx$  is **evaluated** between the **limits**,  $x_i \leq x \leq x_f$
- The answer of a **definite integral** is a **numerical value**
- $\int_{x_i}^{x_f} y(x) dx = [z(x)]_{x_i}^{x_f} = z(x_f) - z(x_i)$
- Example  $y = x \rightarrow \int y dx = \int x dx = \frac{1}{2}x^2 = z(x)$
- $\int_{x_i}^{x_f} x dx = [z(x)]_{x_i}^{x_f} = z(x_f) - z(x_i) = \frac{1}{2}x_f^2 - \frac{1}{2}x_i^2$

# Chain Rule

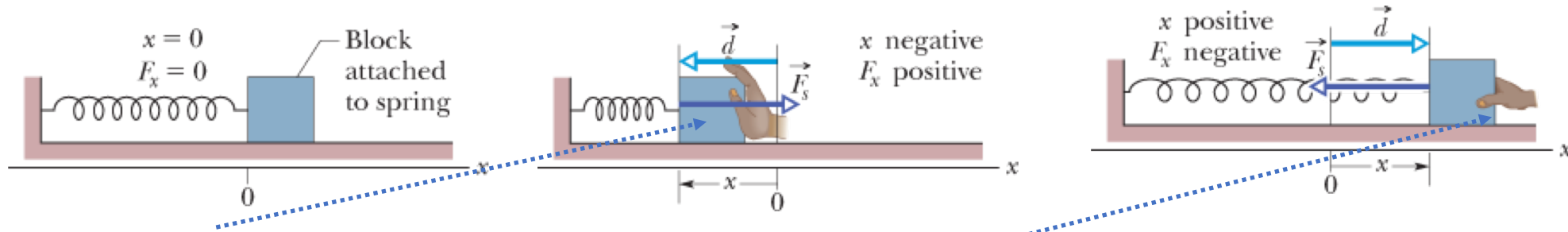
- Let  $x$  be a function of  $t$ ,  $x(t)$
- If  $y$  is a function of  $x$ ,  $y(x)$  then  $y$  is also a function of  $t$ ,  $y(t)$
- Now consider  $\frac{dy}{dt}$ , the **chain rule** states  $\frac{dy}{dt} = \frac{dy}{dx} \frac{dx}{dt}$
- **Example**,  $x(t) = t^2$ ,  $y(x) = x^2 = (t^2)^2 = t^4 = y(t)$
- Since  $y(t) = t^4 \rightarrow \frac{dy}{dt} = 4 \times t^{4-1} = 4t^3$
- Also  $x(t) = t^2 \rightarrow \frac{dx}{dt} = 2 \times t^{2-1} = 2t$
- Noting  $y = x^2$ , use the **chain rule**  $\frac{dy}{dt} = \frac{dy}{dx} \frac{dx}{dt} = \frac{d(x^2)}{dx} \times 2t = 2 \times x^{2-1} \times 2t = 4xt$
- But  $x(t) = t^2$
- This means  $\frac{dy}{dt} = \frac{dy}{dx} \frac{dx}{dt} = 4xt = 4t^3$

# Advance Differential Calculus

- Let  $x$  be a function of  $t$ ,  $x(t)$
- If  $y$  is a function of  $x$ ,  $y(x)$  then  $y$  is also a function of  $t$ ,  $y(t)$
- The **chain rule** states  $\frac{dy}{dt} = \frac{dy}{dx} \frac{dx}{dt}$
- Now consider the **definite integral**  $\int_{x_i}^{x_f} \frac{dy}{dt} dx$
- We can change the variable of integration from  $x$  to ... say  $y$  by doing the following
- $\int_{x_i}^{x_f} y(x) dx = \int_{y_i}^{y_f} \frac{dy}{dx} \frac{dx}{dt} dx$ , where the **integration limits** are now  $y(x_i) = y_i$  and  $y(x_f) = y_f$
- Though I have not proven this,  $\frac{dy}{dx} dx = dy$
- $\int_{x_i}^{x_f} \frac{dy}{dt} dx = \int_{y_i}^{y_f} \frac{dy}{dx} \frac{dx}{dt} dx = \int_{y_i}^{y_f} \frac{dx}{dt} dy$ , where  $\frac{dx}{dt}$  must be expressed as a function of  $y$

# The Spring Force

Figure below shows a spring in an equilibrium (relaxed) state:



If a **leftward** force is applied to the block that **compresses** the spring,  $x < 0$  (**left**), the spring applies a **rightward** restoring spring force,  $\vec{F}_s$ :  $F_x > 0$  pushes on the block.

If a **rightward** force is applied to the block that **stretches** the spring,  $x > 0$  (**right**), the spring applies a **leftward** restoring spring force,  $\vec{F}_s$ :  $F_x < 0$  pulls the block.

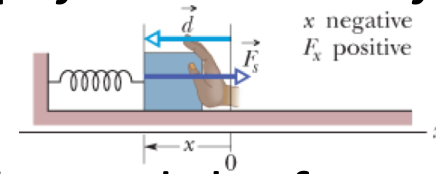
**Hooke's Law:** Spring force  $F_x = -kx$ ,  $x$  is displacement of spring,  $k$  is the spring constant in unit of N/m

# Work by a spring

Work by **variable spring force** in 1D that changes the **compression/stretching** of **spring** from an **initial**  $x_i$  to **final**  $x_f$ .

Divide into **infinitesimal portion**,  $\Delta x_j$ , where the subscript  $j$  denotes the  $j^{\text{th}}$  portion of the path.

Work by spring,  $dW_s = F_{xj} \Delta x_j > 0, F_{xj} < 0, \Delta x_j < 0$



It is assumed that the infinitesimal portion,  $\Delta x_j$ , is straight, and the force is constant,  $F_{xj}$ .

The total work done by spring on the path is found by summing all  $j^{\text{th}}$  portions:

$$W_s = \sum dW = \sum_j F_{xj} \Delta x_j$$

This gives  $W_s = \int_{x_i}^{x_f} F_x(x) dx = - \int_{x_i}^{x_f} kx dx = ?$

What is the definite integral,  $\int kx dx = ?$   $\int kx dx = \frac{1}{2} kx^2$

Definite integral, total work by spring,  $W_s = - \int_{x_i}^{x_f} kx dx = \left[ -\frac{1}{2} kx^2 \right]_{x_i}^{x_f}$

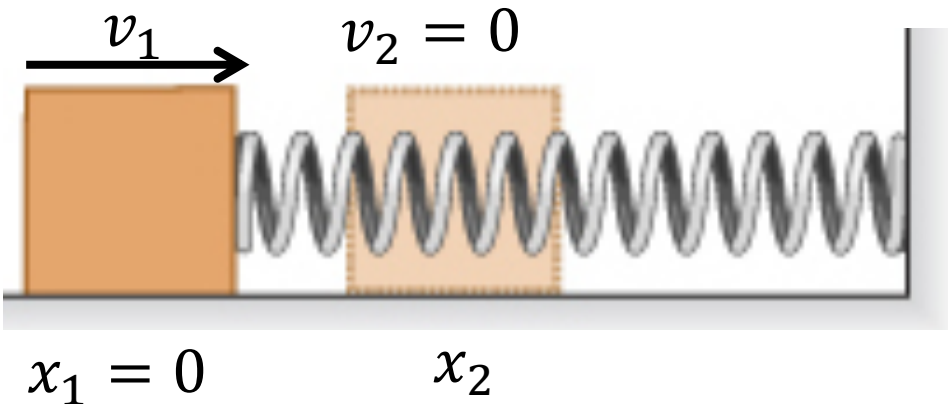
$$W_s = \left( -\frac{1}{2} kx_f^2 \right) - \left( -\frac{1}{2} kx_i^2 \right) = \frac{1}{2} kx_i^2 - \frac{1}{2} kx_f^2$$

# Example spring

A typical spring constant is  $k = 400 \frac{N}{m}$ . Below such a spring is being compressed from **equilibrium**, at  $x_1 = 0$ , by a  $m = 8 \text{ kg}$  box, moving at  $v_1 = 9 \frac{m}{s}$ . Assume no friction.

A) Find maximum compression,  $x_2$ .

B) Find the speed,  $v_3$ , when the box is at the equilibrium at  $x_3 = 0$ .

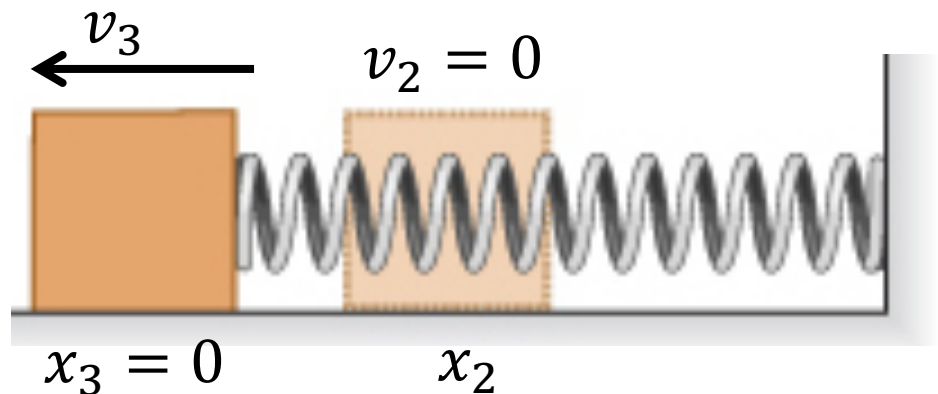


Part A:

$$\text{Work by spring from } x_1 \text{ to } x_2, W_{12} = \frac{1}{2} k x_1^2 - \frac{1}{2} k x_2^2$$

$$\text{Work-Energy, } W_{12} = -\frac{1}{2} k x_2^2 = \Delta K = \frac{1}{2} m v_2^2 - \frac{1}{2} m v_1^2$$

$$x_2 = \sqrt{\frac{m}{k}} v_1 = \sqrt{\frac{8 \text{ kg}}{400 \text{ Nm}^{-1}}} \times 9 \frac{\text{m}}{\text{s}} = 1.27 \text{ m}$$



Part B: Use work-energy theorem from 2 to 3

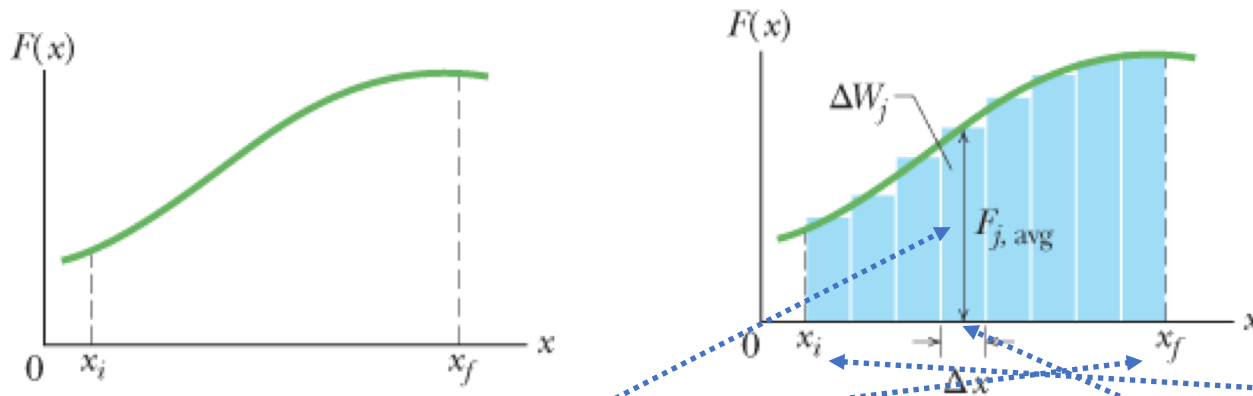
$$W_{23} = \frac{1}{2} k x_2^2 - \frac{1}{2} k x_3^2 = \Delta K = \frac{1}{2} m v_2^2 - \frac{1}{2} m v_3^2$$

$$v_3 = \sqrt{\frac{k}{m}} x_2 = \sqrt{\frac{400 \text{ Nm}^{-1}}{8 \text{ kg}}} \times 1.27 \text{ m} = 8.98 \frac{\text{m}}{\text{s}}$$



# Work by general variable force in 1D

Consider a object moving in 1D acted on by a general variable force as shown below,  $F(x)$ .



The work done is the area below the  $F$  vs.  $x$  curve from the initial position  $x_i$  to final position  $x_f$ .

Divide into very small paths each of length,  $\Delta x$ , where at position  $x_j$  the average force is  $F_{j, \text{avg}}$ , which contributes work  $\Delta W_j = F_{j, \text{avg}} \Delta x$

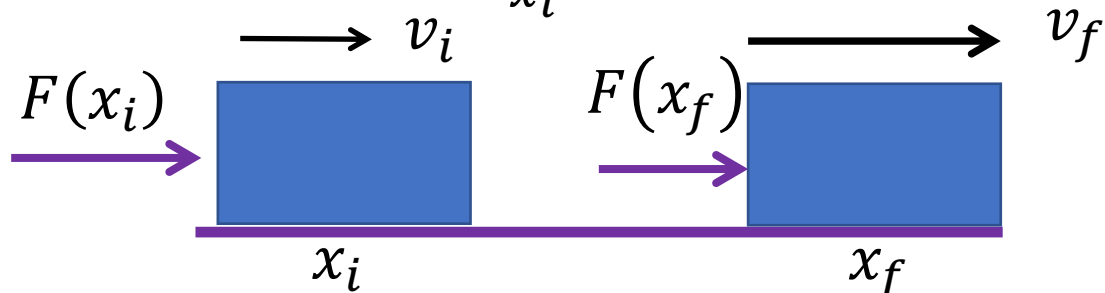
The total work done is  $W \sim \sum_j W_j = \sum_j F_{j, \text{avg}} \Delta x$

In the calculus limit of very small  $\Delta x \rightarrow 0$ , the total work done is

$$W = \int_{x_i}^{x_f} F(x) dx$$

# Work-Energy Theorem by a general variable force in 1D

The work done is  $W = \int_{x_i}^{x_f} F(x) dx$ .



Use Newton's second law,  $F = ma \rightarrow W = \int_{x_i}^{x_f} madx$

Divide Acceleration is  $a = \frac{dv}{dt} \rightarrow adx = \frac{dv}{dt} dx$

Note velocity  $v$  is a function of position  $x$  and time  $t$ .

Use the chain rule,  $\frac{dv}{dt} = \frac{dv}{dx} \frac{dx}{dt} = v \frac{dv}{dx}$ , where  $\frac{dx}{dt} = v$

$$adx = \frac{dv}{dt} dx = v \frac{dv}{dx} dx = v dv$$

$$W = \int_{x_i}^{x_f} madx = m \int_{v_i}^{v_f} v dv = \left[ \frac{1}{2} mv^2 \right]_{v_i}^{v_f} = \frac{1}{2} mv_f^2 - \frac{1}{2} mv_i^2 = K_f - K_i$$

Work done  $W = \Delta K$  change in kinetic energy